

① Let us determine the eigenvalues of A . We have

$$\det(\lambda I - A) = 0 \quad \begin{vmatrix} \lambda-2 & 1 & -2 \\ -1 & \lambda-2 & 4 \\ -3 & 1 & \lambda-1 \end{vmatrix} = 0$$

$$\begin{vmatrix} \lambda+1 & 0 & -(\lambda+1) \\ -1 & \lambda-2 & 4 \\ -3 & 1 & \lambda-1 \end{vmatrix} = 0 \quad \begin{vmatrix} 0 & 0 & -(\lambda+1) \\ +3 & \lambda-2 & 4 \\ \lambda-4 & 1 & \lambda-1 \end{vmatrix} = 0$$

$$-(\lambda+1) [3 - (\lambda-2)(\lambda-4)] = 0$$

$$(\lambda+1)(\lambda^2 - 6\lambda + 8 - 3) = 0$$

$$(\lambda+1)(\lambda^2 - 6\lambda + 5) = 0$$

$$\lambda_1 = 1$$

$$\lambda_2 = -1$$

$$\lambda_3 = 5$$

Let us now evaluate the corresponding eigenvectors. We have

$$\begin{bmatrix} -1 & 1 & -2 \\ -1 & -1 & 4 \\ -3 & 1 & 0 \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \\ \gamma \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \quad \begin{cases} -\alpha - \beta + 4\gamma = 0 \\ -3\alpha + \beta = 0 \end{cases}$$

$$\vec{h}_1 = \begin{bmatrix} L \\ 3L \\ L \end{bmatrix}$$

$$\begin{bmatrix} -3 & 1 & -2 \\ -1 & -3 & 4 \\ -3 & 1 & -2 \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \\ \gamma \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \quad \begin{cases} -6\alpha + 2\beta - 4\gamma = 0 \\ -\alpha - 3\beta + 4\gamma = 0 \end{cases}$$

$$\vec{h}_2 = \begin{bmatrix} M \\ -7M \\ -5M \end{bmatrix}$$

$$\begin{bmatrix} 3 & 1 & -2 \\ -1 & 3 & 4 \\ -3 & 1 & 4 \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \\ \gamma \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \quad \begin{cases} 2\beta + 2\gamma = 0 \\ -\alpha + 3\beta + 4\gamma = 0 \\ -3\alpha + \beta + 4\gamma = 0 \end{cases}$$

$$\underline{h}_3 = \begin{bmatrix} -N \\ N \\ -N \end{bmatrix}$$

Therefore

$$\underline{z} = \begin{bmatrix} e^x & e^{-x} & -e^{5x} \\ 3e^x & -7e^{-x} & e^{5x} \\ e^x & -5e^{-x} & -e^{5x} \end{bmatrix} \begin{bmatrix} L \\ M \\ N \end{bmatrix}$$

$$\textcircled{2} \begin{cases} y' = \frac{xy}{1-xy} \\ y(x_0) = y_0 \end{cases}$$

$$D = \{(x, y) \in \mathbb{R}^2 : xy \neq 1\}$$

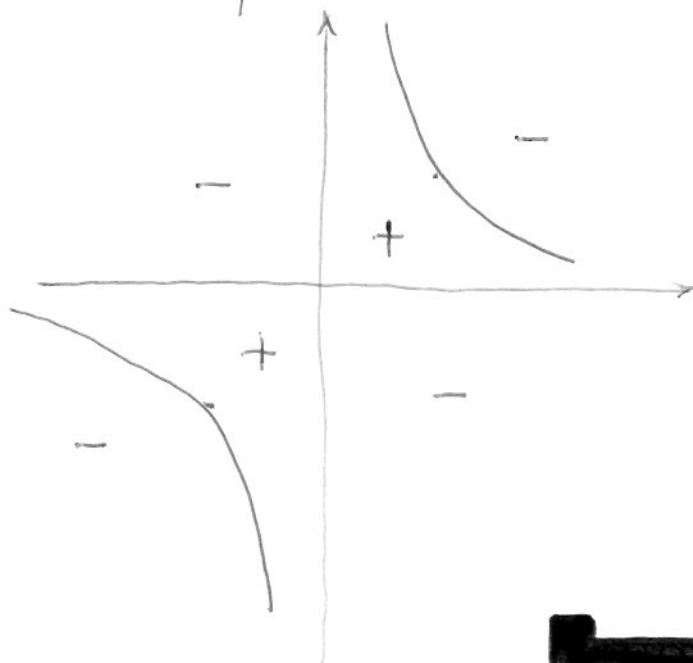
D is an open set and it is apparent that $f \in C^\infty(D)$

Therefore, we can apply both the local existence and uniqueness Theorem and the regularity theorem to conclude that

$\forall (x_0, y_0) \in D \quad \exists! y = y(x) \quad y \in C^\infty(I_{x_0})$
solution to the Cauchy problem stated above.

The global existence and uniqueness theorem cannot be applied.

The figure on the right describes where the function is increasing (+) and where it is decreasing (-).



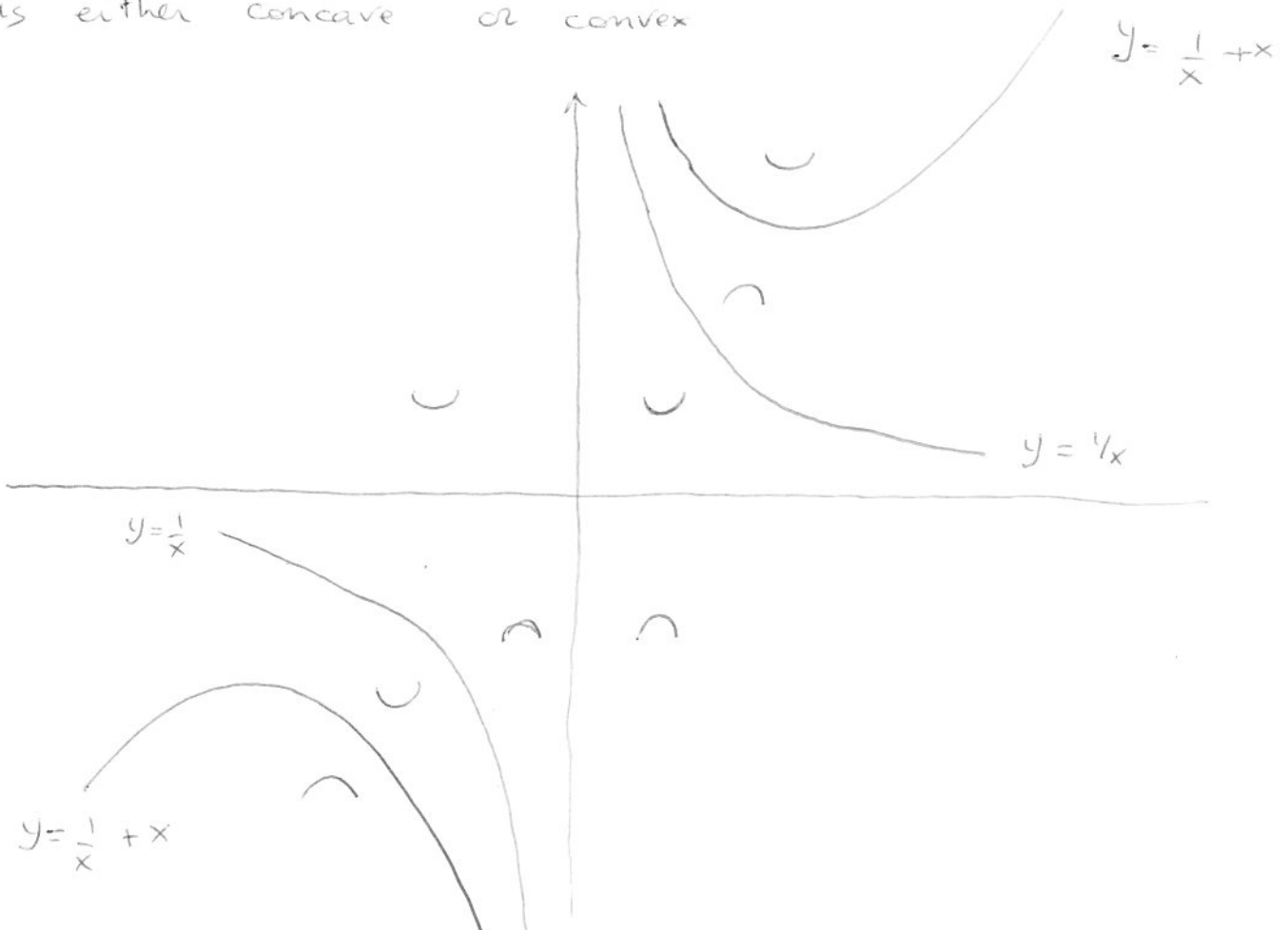
Notice that if (x_1, y_1) is such that $x_1 y_1 = 1$, then

$$\lim_{x \rightarrow x_1} y'(x) = \infty.$$

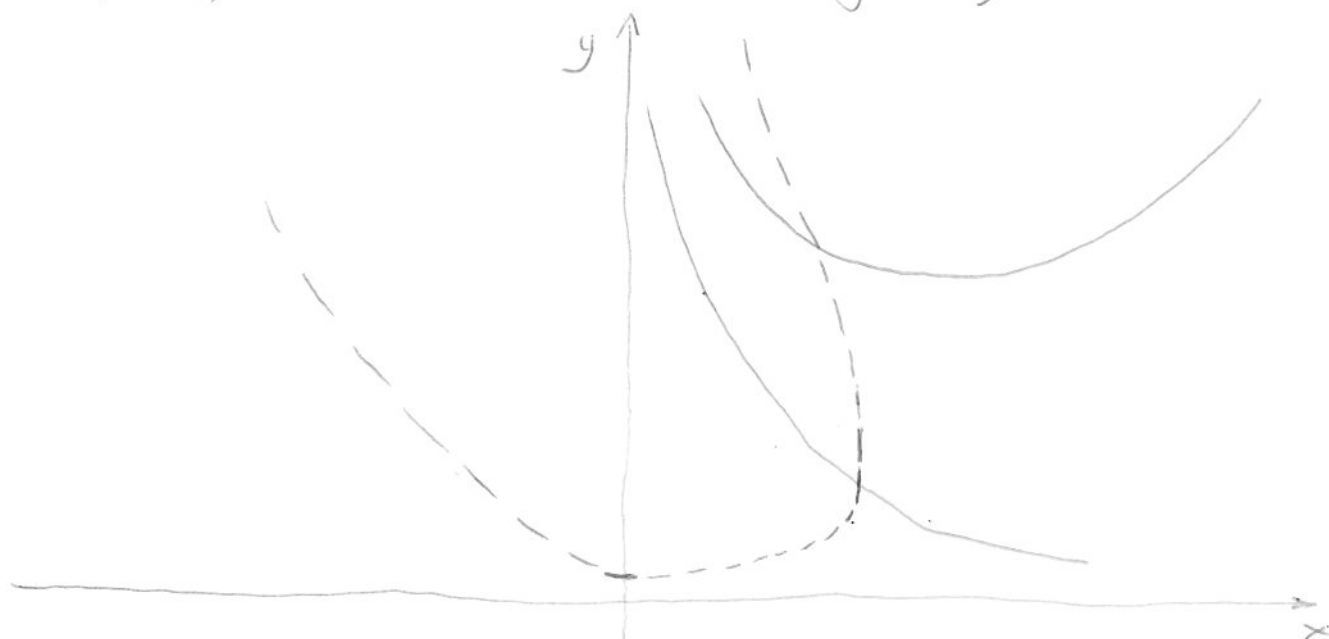
Coming to the second order derivatives, we have

$$\begin{aligned} y'' &= \frac{(y + xy')(1 - xy) + xy(y + xy')}{(1 - xy)^2} = \frac{y + xy'}{(1 - xy)^2} \\ &= \frac{y + x \frac{xy}{1 - xy}}{(1 - xy)^2} = \frac{y - xy^2 + x^2 y}{(1 - xy)^3} \\ &= \frac{y[1 + x^2 - xy]}{(1 - xy)^2(1 - xy)} \end{aligned}$$

Therefore, the figure below gives the regions where y is either concave or convex



Taking into account all the previous information, a qualitative graph has the following behavior (for the sake of simplicity, we restrict ourselves to $y > 0$)



③ $u \in \mathcal{S}'(\mathbb{R})$ since it is a regular function whose growth at infinity is controlled by a polynomial of arbitrary degree.

We now have

$$u' = \frac{2x}{1+x^2}$$

If we transform both sides and take into account the fundamental properties of the Fourier transform, we have

$$\begin{aligned} i \xi \hat{u} &= 2i \frac{d}{d\xi} \pi e^{-|\xi|} \\ \xi \hat{u} &= -2\pi e^{-|\xi|} \frac{|\xi|}{\xi} \\ \xi \hat{u} &= -2\pi e^{-|\xi|} \frac{\xi}{|\xi|} \end{aligned}$$

If we proceed formally, simplifying both sides, we can conclude that

$$\hat{u} = -2\pi \text{ pf } \frac{e^{-|\xi|}}{|\xi|}$$

where

$$\text{pf } \frac{e^{-|\xi|}}{|\xi|} = \lim_{\epsilon \rightarrow 0} \frac{e^{-|\xi|}}{|\xi|} \chi_{(-\infty, -\epsilon) \cup (\epsilon, +\infty)} + 2 \ln \epsilon \delta$$

where the limit is taken in the sense of \mathcal{S}' . Let us show that such a limit exists. By definition, $\forall \varphi \in \mathcal{S}$ we have

$$\langle \text{pf } \frac{e^{-|\xi|}}{|\xi|}, \varphi \rangle = \lim_{\epsilon \rightarrow 0} \int_{-\infty}^{-\epsilon} \frac{e^{\xi}}{\xi} \varphi(\xi) + \int_{\epsilon}^{+\infty} \frac{e^{-\xi}}{\xi} \varphi(\xi) + 2 \ln \epsilon \varphi(0)$$

We consider

$$\lim_{\epsilon \rightarrow 0} \int_{\epsilon}^{+\infty} \frac{e^{-\xi}}{\xi} \varphi(\xi) + \ln \epsilon \varphi(0)$$

since the other term is analogous. We can rewrite

$$\varphi(\xi) = \varphi(0) + \xi \psi(\xi)$$

where $\psi \in \mathcal{S}$. Therefore

$$\begin{aligned} \int_{\epsilon}^{+\infty} \frac{e^{-\xi}}{\xi} \varphi(\xi) + \ln \epsilon \varphi(0) &= \int_{\epsilon}^{\infty} \frac{e^{-\xi}}{\xi} (\varphi(0) + \xi \psi(\xi)) + \ln \epsilon \varphi(0) \\ &= \int_{\epsilon}^{\infty} \frac{e^{-\xi}}{\xi} \varphi(0) + \ln \epsilon \varphi(0) + \int_{\epsilon}^{\infty} e^{-\xi} \psi(\xi) \end{aligned}$$

When $\epsilon \rightarrow 0$, the last term gives no problem. Let us concentrate on the first two ones. If we integrate by parts, we have

$$\begin{aligned} \int_{\epsilon}^{\infty} \frac{e^{-\xi}}{\xi} \varphi(0) + \ln \epsilon \varphi(0) &= e^{-\xi} \ln \xi \Big|_{\epsilon}^{\infty} \varphi(0) + \int_{\epsilon}^{\infty} e^{-\xi} \ln \xi \\ \varphi(0) + \ln \epsilon \varphi(0) &= -e^{-\epsilon} \ln \epsilon \varphi(0) + \ln \epsilon \varphi(0) + \varphi(0) \int_{\epsilon}^{\infty} e^{-\xi} \ln \xi \\ &= \epsilon \varphi(0) \ln \epsilon \frac{1-e^{-\epsilon}}{\epsilon} + \varphi(0) \int_{\epsilon}^{\infty} e^{-\xi} \ln \xi \end{aligned}$$

When $\epsilon \rightarrow 0$ the first term vanishes, whereas the second one has finite limit.

(4) From v_0 we build u_0 .

We have

$$u_0 = \frac{v_0}{\|v_0\|_{H^1}} = \frac{1}{\left[\int_{-1}^1 |v_0|^2 dx + \int_{-1}^1 |v_0'|^2 dx \right]^{1/2}} = \frac{1}{\sqrt{2}}$$

Let us now build u_1

$$u_1 = \frac{v_1 - (v_1, u_0)u_0}{\|v_1 - (v_1, u_0)u_0\|_{H^1}}$$

$$\begin{aligned} (v_1, u_0) &= \int_{-1}^1 v_1 u_0 dx + \int_{-1}^1 v_1' u_0' dx = \int_{-1}^1 \frac{1}{\sqrt{2}} x dx \\ &= \frac{1}{\sqrt{2}} \frac{x^2}{2} \Big|_{-1}^1 = \frac{1}{\sqrt{2}} \cdot 0 = \cancel{0} \end{aligned}$$

$$v_1 - (v_1, u_0)u_0 = x$$

$$\begin{aligned} \|v_1 - (v_1, u_0)u_0\|_{H^1}^2 &= \int_{-1}^1 x^2 dx + \int_{-1}^1 1^2 dx \\ &= 2 \frac{1}{3} + 2 = \frac{8}{3} \end{aligned}$$

Quindi

$$u_1 = \frac{x}{\sqrt{\frac{8}{3}}} = \sqrt{\frac{3}{8}} x$$

Infine

$$u_2 = \frac{v_2 - (v_2, u_0)u_0 - (v_2, u_1)u_1}{\|v_2 - (v_2, u_0)u_0 - (v_2, u_1)u_1\|}$$

We have

$$(v_2, u_0) = \int_{-1}^1 x^2 \frac{1}{\sqrt{2}} + \int_{-1}^1 2x \cdot 0 = \frac{2}{\sqrt{2}} \frac{1}{3} = \frac{\sqrt{2}}{3}$$

$$(v_2, u_1) = \int_{-1}^1 x^2 \cdot \sqrt{\frac{3}{8}} x + \int_{-1}^1 2x \cdot \sqrt{\frac{3}{8}} = 0.$$

$$v_2 - (v_2, u_0)u_0 - (v_2, u_1)u_1 = x^2 - \frac{\sqrt{2}}{3} \cdot \frac{1}{\sqrt{2}} = \frac{3x^2-1}{3}$$

$$\|v_2 - (v_2, u_0)u_0 - (v_2, u_1)u_1\|^2 = \int_{-1}^1 \left(\frac{3x^2-1}{3}\right)^2 dx + \int_{-1}^1 (2x)^2 dx$$

$$= \frac{1}{9} \cdot 2 \cdot \int_0^1 (9x^4 - 6x^2 + 1) dx + 2 \cdot 4 \left[\frac{x^3}{3} \right]_0^1 =$$

$$= \frac{2}{9} \left(\cancel{9} \cdot \frac{1}{\cancel{9}} - \frac{2}{\cancel{2}} + 1 \right) + \frac{8}{3} = \frac{2}{9} \left(\frac{18 - 20 + 10}{10} \right) + \frac{8}{3}$$

$$= \frac{2 \cdot \cancel{8}}{90} + \frac{240}{90} = \frac{256}{90} = \frac{128}{45}$$

Finally

$$u_2 = \frac{3x^2-1}{3} / \sqrt{\frac{128}{45}} = \frac{3x^2-1}{3} \frac{3\sqrt{5}}{\sqrt{128}}$$

$$= \sqrt{\frac{5}{128}} (3x^2-1).$$